Local resonances and parametric level dynamics in the many-body localised phase arXiv:2107.12387

Sam Garratt with Sthitadhi Roy and John Chalker

MBL2021: dead or alive?









▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQで

Many-body localised phase

Theory based on existence of local integrals of motion (LIOM).

Serbyn, Papic, Abanin PRL 2013, Huse, Nandkishore, Oganesyan PRB 2014

Example: disordered Heisenberg chain

$$H = \sum_{j} ec{h}_{j} \cdot ec{\sigma}_{j} + J \sum_{j} ec{\sigma}_{j} \cdot ec{\sigma}_{j+1}$$

Many-body localised phase

Theory based on existence of local integrals of motion (LIOM). Serbyn, Papic, Abanin PRL 2013, Huse, Nandkishore, Oganesyan PRB 2014 Example: disordered Heisenberg chain

$$H = \sum_{j} \vec{h}_{j} \cdot \vec{\sigma}_{j} + J \sum_{j} \vec{\sigma}_{j} \cdot \vec{\sigma}_{j+1}$$

For J = 0, LIOM $\tau_j^z = (\vec{h}_j/h_j) \cdot \vec{\sigma}_j$.

Many-body localised phase

Theory based on existence of local integrals of motion (LIOM). Serbyn, Papic, Abanin PRL 2013, Huse, Nandkishore, Oganesyan PRB 2014 Example: disordered Heisenberg chain

$$H = \sum_{j} \vec{h}_{j} \cdot \vec{\sigma}_{j} + J \sum_{j} \vec{\sigma}_{j} \cdot \vec{\sigma}_{j+1}$$

A D N A 目 N A E N A E N A B N A C N

For J = 0, LIOM $\tau_j^z = (\vec{h}_j/h_j) \cdot \vec{\sigma}_j$.

For small $J \neq 0$, LIOM $\tilde{ au}_j^z = au_j^z +$ exponential tails

Resonances

In Anderson localised phase



▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

Resonances

In Anderson localised phase



(日) (四) (日) (日) (日)

Analogous phenomena in MBL phase

- Gopalakrishnan et al. PRB 92 104202 (2015)
- Imbrie PRL 117 027201 (2016)
- Crowley & Chandran arXiv:2012.14393
- Morningstar et al. arXiv:2107.05642

Vary disorder realisation to induce resonances. Avoided level crossings = resonances

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

Vary disorder realisation to induce resonances. Avoided level crossings = resonances

Develop theory in terms of properties of 'standard' LIOM

▲□▶ ▲□▶ ▲ □▶ ▲ □▶ □ のへぐ

Vary disorder realisation to induce resonances. Avoided level crossings = resonances

- Resonances as features in spectra of **local** evolution operators.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Vary disorder realisation to induce resonances. Avoided level crossings = resonances

- Resonances as features in spectra of **local** evolution operators.

・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・

▶ 'Locally pairwise'; eigenstates participate in ~ *L* resonances.

Vary disorder realisation to induce resonances. Avoided level crossings = resonances

- Resonances as features in spectra of local evolution operators.
- ▶ 'Locally pairwise'; eigenstates participate in ~ *L* resonances.
- ► Find small J ≠ 0 distinct from J = 0 for e.g. Distributions of matrix elements of local operators.
- Capture behaviour not usually considered 'localised'.

Avoided crossings as resonances



From here: Floquet analogue of disordered Heisenberg chain

(日) (四) (日) (日) (日)

Avoided crossings as resonances



From here: **Floquet** analogue of disordered Heisenberg chain Floquet operator W, with $W | n \rangle = e^{i\theta_n} | n \rangle$, $n = 1 \dots 2^L$ θ_n = quasienergy, average d.o.s. $\sum_n \langle \delta_{2\pi}(\theta - \theta_n) \rangle = [2\pi]^{-1} \times 2^L$.

Level curvatures

Distribution of level curvatures $\kappa_n = \partial_\lambda^2 \theta_n$. At large κ

$$p_\kappa(\kappa) \sim L |\kappa|^{-(2-\zeta \ln 2)}$$

 $\zeta={\rm decay}$ length for standard LIOM not involved in resonances.

Level curvatures

Distribution of level curvatures $\kappa_n = \partial_\lambda^2 \theta_n$. At large κ

$$p_\kappa(\kappa) \sim L |\kappa|^{-(2-\zeta \ln 2)}$$

 $\zeta = {\rm decay}$ length for standard LIOM not involved in resonances.



Local observables

LIOM at J = 0 are single-site $\tau_j^z = (\vec{h}_j/h_j) \cdot \vec{\sigma}_j$. \implies For J = 0, $\langle n | \tau_j^z | m \rangle = 0$ for $n \neq m$.

Local observables

LIOM at J = 0 are single-site $\tau_j^z = (\vec{h}_j/h_j) \cdot \vec{\sigma}_j$. \implies For J = 0, $\langle n | \tau_j^z | m \rangle = 0$ for $n \neq m$. For $J \neq 0$? Not so!

$$Z = |\langle n | \tau_j^z | m \rangle|^2$$
$$p_Z(Z) \sim 2^{-L} Z^{-3/2}$$

•

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Parametric approach gives behaviour on/off resonance \rightarrow distributions.

Local observables

LIOM at J = 0 are single-site $\tau_j^z = (\vec{h}_j/h_j) \cdot \vec{\sigma}_j$. \implies For J = 0, $\langle n | \tau_j^z | m \rangle = 0$ for $n \neq m$. For $J \neq 0$? Not so!

$$Z = |\langle n | \tau_j^z | m \rangle|^2$$
$$p_Z(Z) \sim 2^{-L} Z^{-3/2}$$

Parametric approach gives behaviour on/off resonance \rightarrow distributions.



(日) (四) (日) (日) (日)



(ロ)、(型)、(E)、(E)、(E)、(O)へ(C)

- Effective $\zeta = \zeta(J)$ is an average over space and disorder.
- ▶ From perturbation theory expect $e^{-p/\zeta} \sim J^p$, or

$$\zeta(J) = \frac{1}{\ln[J_0/J]}$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

- Effective $\zeta = \zeta(J)$ is an average over space and disorder.
- From perturbation theory expect $e^{-p/\zeta} \sim J^p$, or

$$\zeta(J) = \frac{1}{\ln[J_0/J]}$$

From theory + numerics on physical quantities, infer ζ



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ ○臣 - の々ぐ

Vary disorder realisation to induce resonances. Avoided level crossings = resonances

- Resonances as features in spectra of local evolution operators.
- ▶ 'Locally pairwise'; eigenstates participate in ~ *L* resonances.
- ► Find small J ≠ 0 distinct from J = 0 for e.g. Distributions of matrix elements of local operators.
- Capture behaviour not usually considered 'localised'.

Spectral function see also Gopalakrishnan et al. PRB 2015, Crowley & Chandran arXiv:2012.14393



◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Spectral statistics

Two-point correlator of the level density

$$p_{\omega}(\omega) = [2\pi]^{-1} \left[1 - a \frac{L}{2^L} \omega^{-\zeta \ln 2} + \dots \right]$$



~ ~ ~ ~